

Flavoured Co-annihilations

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D. Choudhury, R. Garani, S. Vempati [arXiv:1104.4467](https://arxiv.org/abs/1104.4467) [hep-ph]

Neutralino Dark Matter

Supersymmetric theories have a few dark matter candidates depending on model of SUSY breaking

Neutralino, Gravitino, axino, saxino, etc..

The most popular of them is the minimal Supergravity/CMSSM model with neutralino as the LSP

The model has just four[±] parameters which are
 m_0 $M_{1/2}$ $\tan\beta$ $\text{sg}(\mu)$ A_0

The entire mass spectrum and couplings are determined by these parameters

Neutralino Mass matrix

$$\mathcal{L} \supset \frac{1}{2} \Psi_N \mathcal{M}_N \Psi_N^T + H.c \quad \Psi_N = \{\tilde{B}, \tilde{W}^0, \tilde{H}_1^0, \tilde{H}_2^0\}$$

$$\mathcal{M}_N = \begin{pmatrix} M_1 & 0 & -M_Z c\beta s\theta_W & M_Z s\beta s\theta_W \\ 0 & M_2 & M_Z c\beta c\theta_W & M_Z s\beta c\theta_W \\ -M_Z c\beta s\theta_W & M_Z c\beta c\theta_W & 0 & -\mu \\ M_Z s\beta s\theta_W & -M_Z s\beta c\theta_W & -\mu & 0 \end{pmatrix}$$

Diagonalising

$$N^* \cdot M_{\tilde{N}} \cdot N^\dagger = \text{Diag.}(m_{\chi_1^0}, m_{\chi_2^0}, m_{\chi_3^0}, m_{\chi_4^0})$$

$$\chi_1^0 = N_{\tilde{B}1} \tilde{B} + N_{\tilde{W}1} \tilde{W}^0 + N_{\tilde{H}_11} \tilde{H}_1^0 + N_{\tilde{H}_21} \tilde{H}_2^0$$

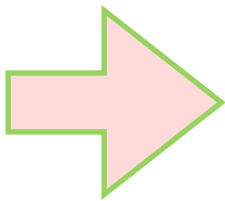
The composition of the lightest neutralino determines the regions where the relic density is satisfied

Three main regions

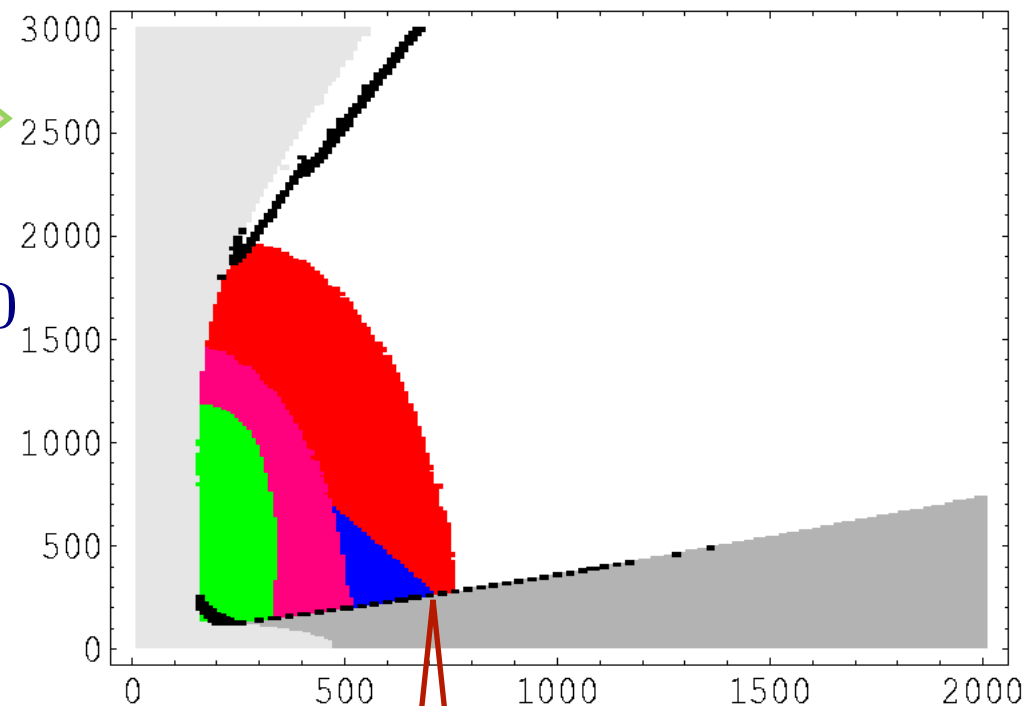
Focus Point/HP Region

$$m_0 \gg M_{1/2}$$

significant higgsino component



m_0



Djouadi, Drees
and Kneur, 2006

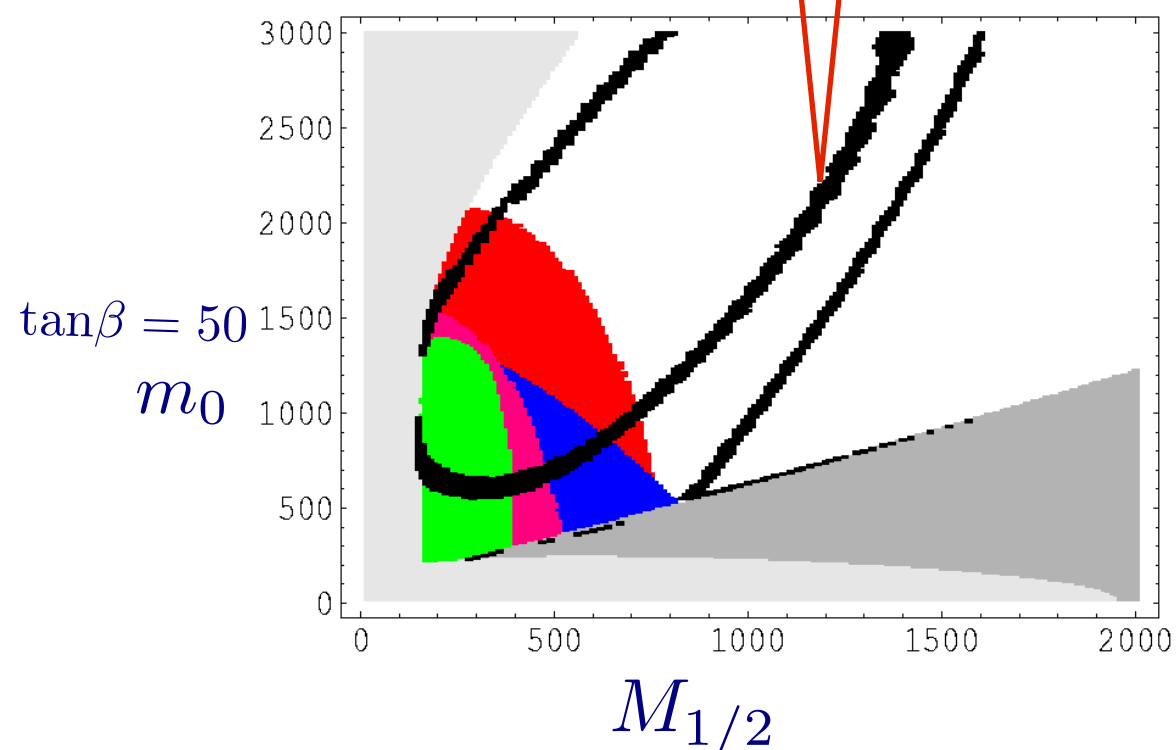
$$\tan\beta = 30$$

$$A_0 = 0$$

$$m_t = 173.2\text{GeV}$$

Funnel Region

$$2m_{\chi_0^1} = m_A$$



Stau-Coannihilation Region

$$m_{\chi_0^1} \approx m_{\tilde{\tau}_R}$$

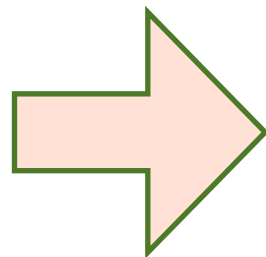
I have not discussed revival of bulk and stop
co-annihilation regions with large A-terms

Completely different phenomenology

BUT TWO POINTS

The scale of mediation of supersymmetry breaking is close to Planck Scale. The GUT scale is at least two orders of magnitude lower than that.

$$M_{\text{Pl}} = 10^{18} \text{ GeV}$$



$$M_{\text{GUT}} = 2 \cdot 10^{16} \text{ GeV}$$

*It is well known that RG running between these two scales can significantly modify the low energy spectrum, especially the flavour sector

Hall, Kostelcky, Raby, 87
Barbieri, Hall, Strumia, 96

*Neutrino masses are well established now. Seesaw mechanism needs to be incorporated. Fits naturally in SO(10) models. Has again, large effects in weak scale physics in terms of flavour violation.

Masiero Borzumati 86
Casas Ibarra 01

No reason to expect that these effects would not modify the dark matter regions.

GUT Effects on Dark Matter

(1) Through flavour violating entries in the slepton mass matrix with and without seesaw
(Modifies Co-annihilation regions)

not explored; topic of the present talk

(2) RG corrections to the Higgs sector
especially in the presence of a Seesaw
(Significantly modifies the focus point regions)

(3) Through gauge corrections to the sleptonic fields
(as they sit in the same multiplet as squarks)
(modifies Co-annihilation regions)

Calibbi et. al, '06, '07, '10 Goto et. al, '07

Barger et.al, '08, '10 Gomez et.al, '08, '09, '10

Olive et.al, '08, '09, '10

The effect of LFV

- (a) The mass of the lightest stau gets modified
- (b) The cross sections get modified in the early universe because of the presence of flavour violation

Both these effects significantly modify the stau co-annihilation region

Flavour Effects in the co-annihilation region

In GUT theories like SU(5),

$$W = h_{ij}^u 10_i 10_j \bar{5}_H + h_{ij}^d 10_i \bar{5}_j 5_H + \dots$$

RG running above the Grand Unified Scale Can give rise to Flavour Violation in the RR sector

$$\begin{aligned} (\Delta_{RR})_{i \neq j} &= (m_{\tilde{10}}^2)_{ij} (M_X \rightarrow M_{GUT}) \\ &= -3 \cdot \frac{3m_0^2 + a_0^2}{16\pi^2} V_{ti} V_{tj} \ln \left(\frac{M_X^2}{M_{GUT}^2} \right) h_t^2 \end{aligned}$$

Suppressed by CKM angles.

Limits from Flavour Violation much weaker

Left-Right Symmetric Models

Left Right symmetric models could have sizeable
LFV in the RR sector

*C. S. Aulakh et. al
'97 '98*

$$\Delta m_L^2 = -\frac{1}{4\pi^2} (3f_c f_c^\dagger + Y_L Y_L^\dagger) (3m_0^2 + A_0^2) \ln \left(\frac{m_{\text{GUT}}}{v_r} \right)$$

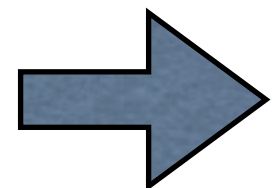
$$\Delta m_{L^c}^2 = -\frac{1}{4\pi^2} (3f_c^\dagger f_c + Y_L^\dagger Y_L) (3m_0^2 + A_0^2) \ln \left(\frac{m_{\text{GUT}}}{v_r} \right)$$

*Esteves et. al
arXiv: 1011.0348*

However, LR symmetry forces same amount of LFV in
the LL sector too !

$\delta \sim \mathcal{O}(10^{-1})$ could still be generated

Model Independent Approach



Assume $(\Delta)_{RR}$

(FN Models, moduli breaking etc)

The weak scale mass matrix $\Phi^T \mathcal{M}_{\tilde{l}}^2 \Phi$ in the basis $\Phi^T = \{\tilde{\mu}_L, \tilde{\tau}_L, \tilde{\mu}_R, \tilde{\tau}_R\}$ has the form :

$$\mathcal{M}_{\tilde{l}}^2 = \begin{pmatrix} m_{\tilde{\mu}_L}^2 & 0 & m_{\tilde{\mu}_{LR}}^2 & 0 \\ 0 & m_{\tilde{\tau}_L}^2 & 0 & m_{\tilde{\tau}_{LR}}^2 \\ m_{\tilde{\mu}_{LR}}^2 & 0 & m_{\tilde{\mu}_R}^2 & \Delta_{RR}^{\mu\tau} \\ 0 & m_{\tilde{\tau}_{LR}}^2 & \Delta_{RR}^{\mu\tau} & m_{\tilde{\tau}_R}^2 \end{pmatrix}$$

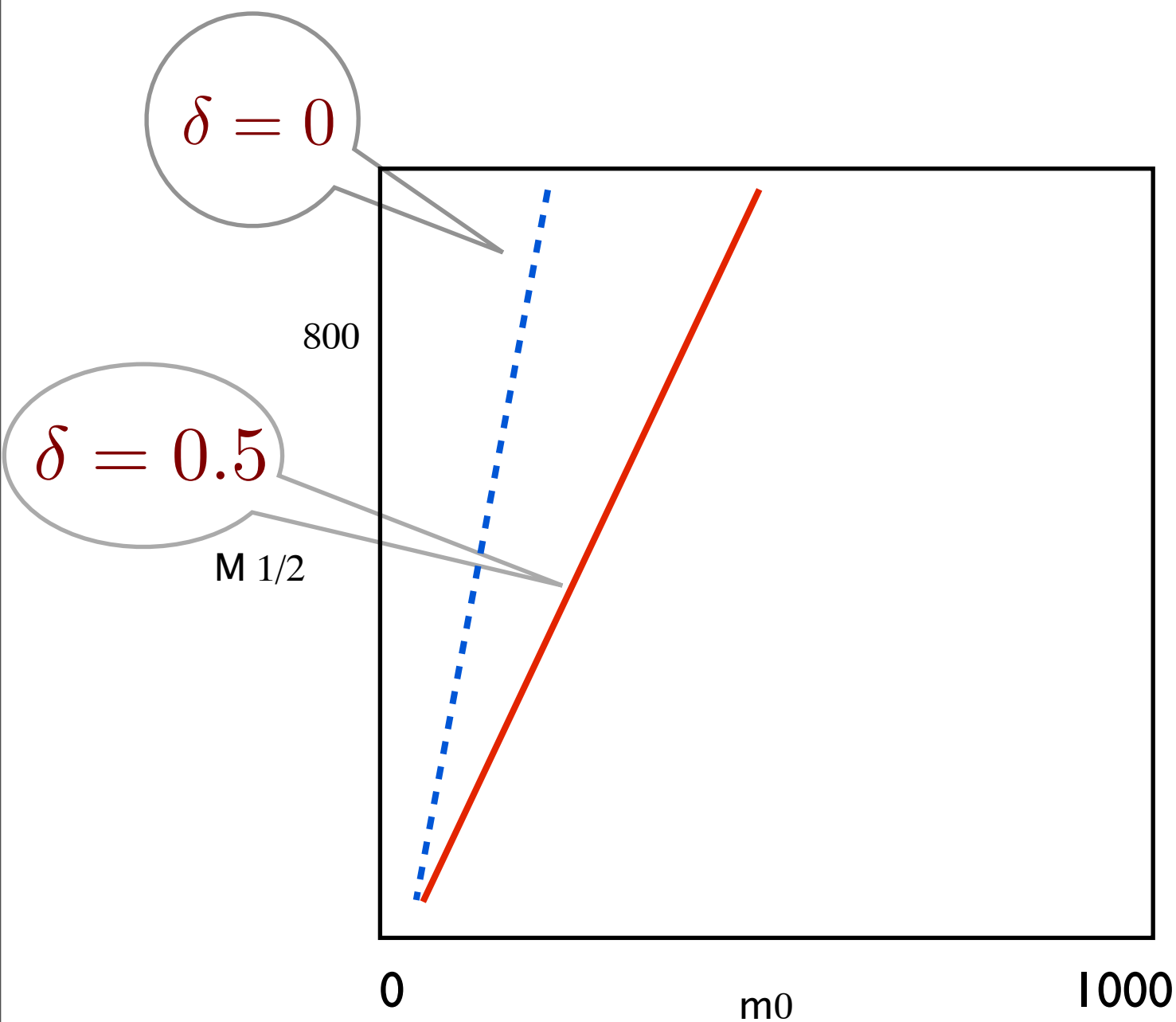
The lightest eigenvalue takes the form :

$$m_{\tilde{\tau}_1}^2 \approx m_{\tilde{\tau}_R}^2 (1 - \delta) - m_\tau \mu \tan \beta$$

No tachyons

where $\delta = (\Delta_{RR})/m_{\tilde{\tau}_R}^2$

$$\delta < 1 - \frac{m_\tau \mu \tan \beta}{m_{\tilde{\tau}_R}^2}$$



The co-annihilation condition is given by

$$m_{\tilde{\tau}_1} \approx M_{\chi_0^1}$$

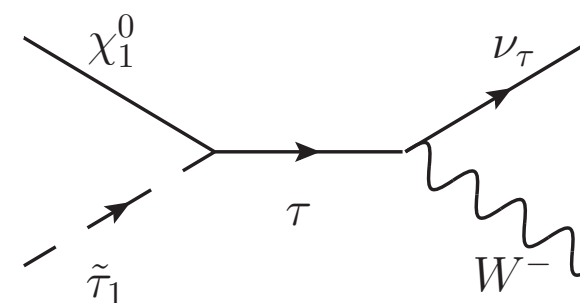
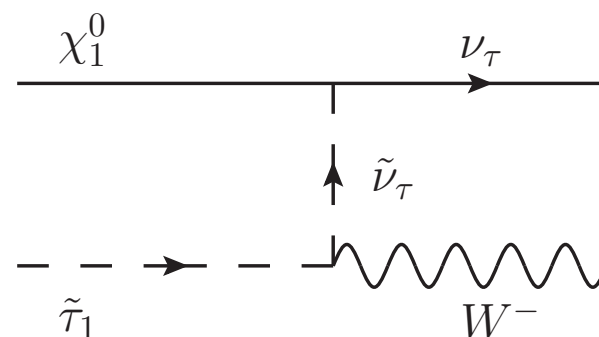
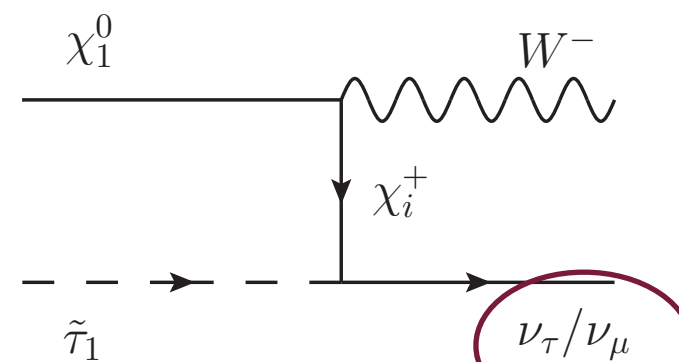
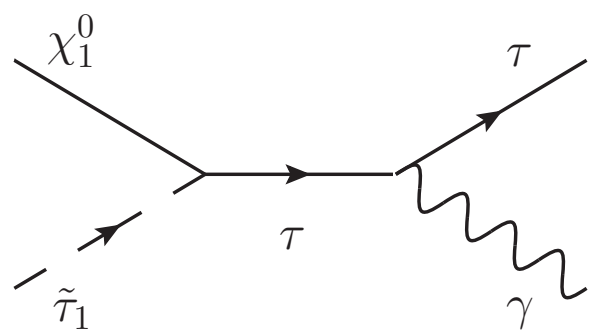
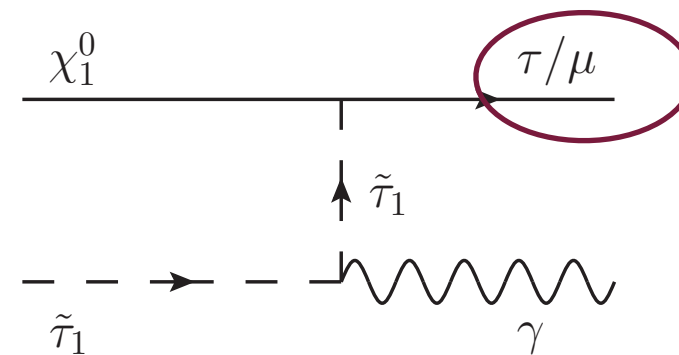
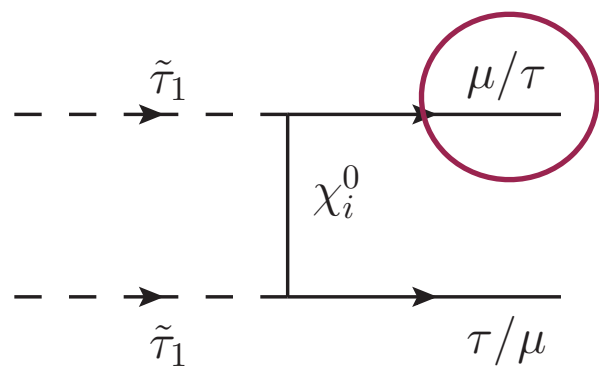
In terms of high energy parameters

Flavour Violation changes the slope, lowers the neutralino mass consistent with relic density

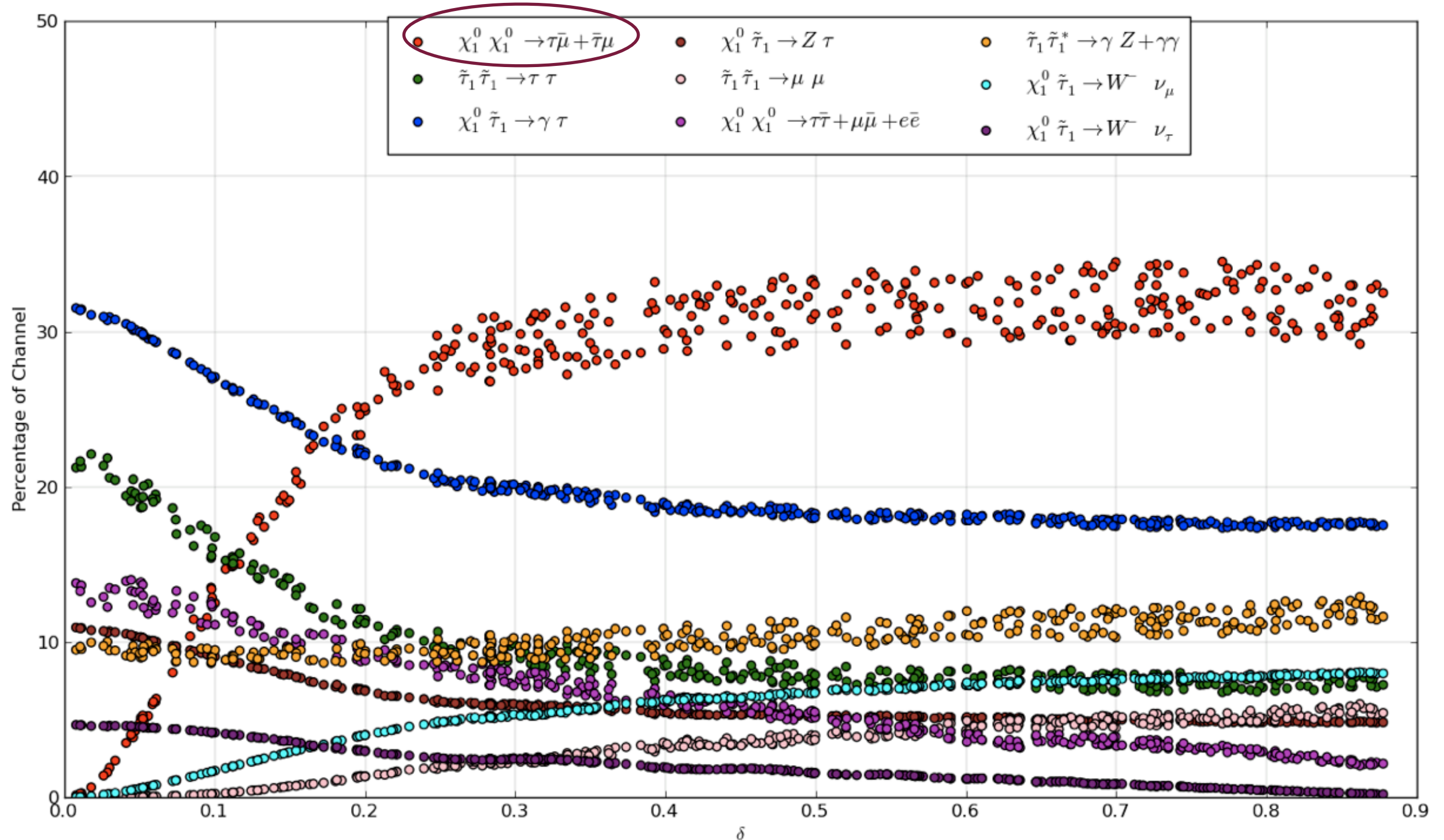
Neglecting flavour (and other) constraints for the time being..

Stau Co-annihilations in the presence of flavour violation

New Flavour violating processes could contribute to the scatterings in the early universe



Channels in Co-annihilation regions



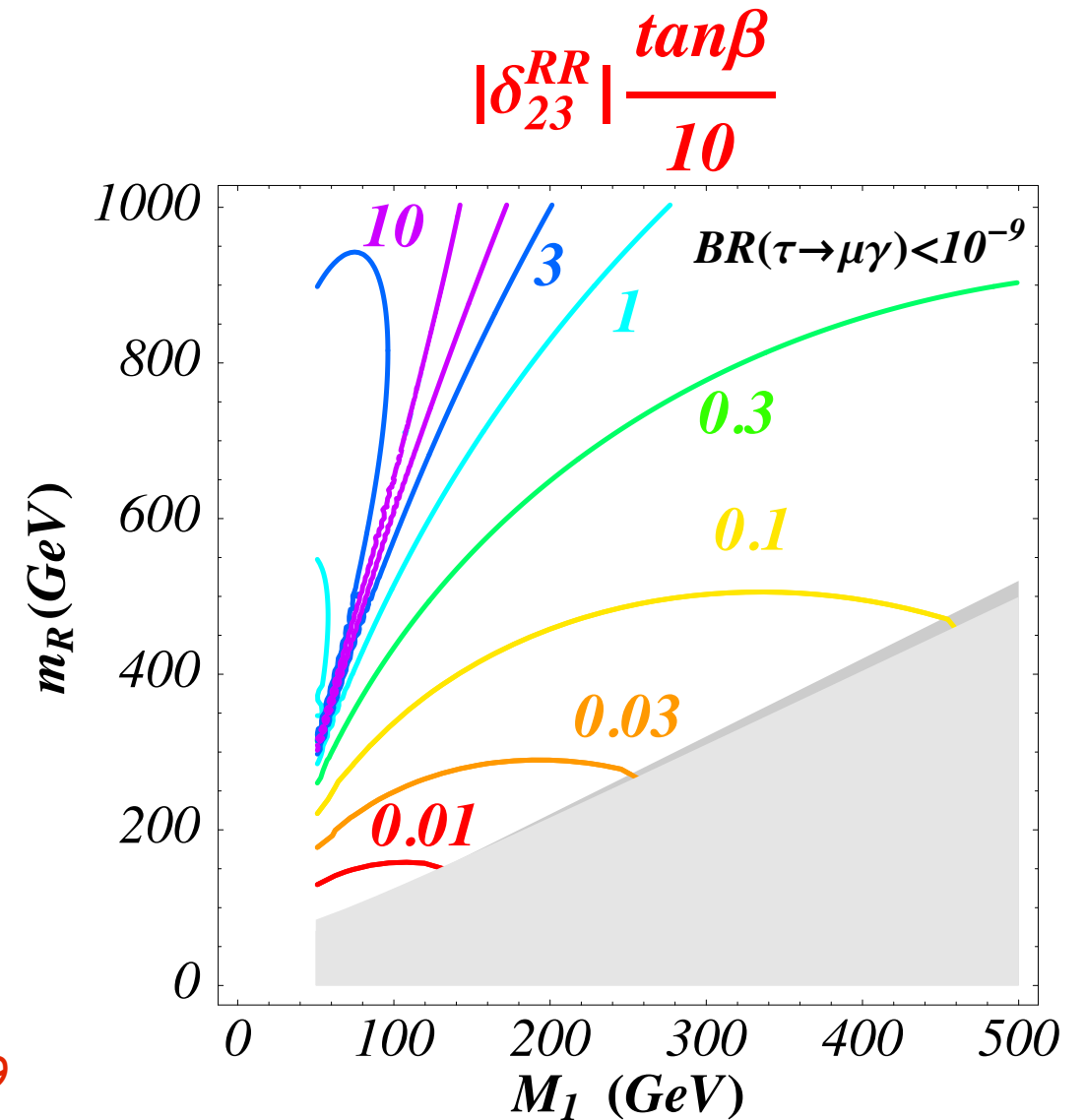
Flavour violating channels become dominant

Choudhury, Garani, Vempati / 104.4467

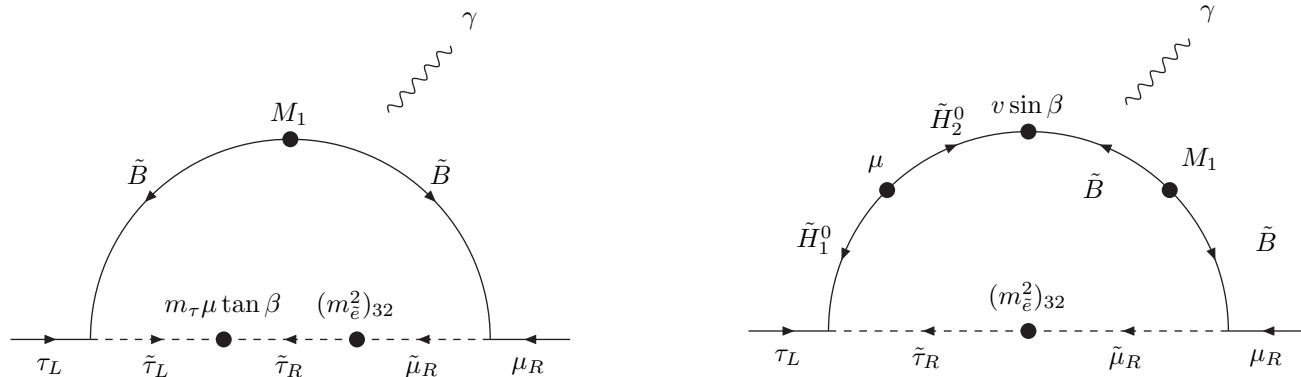
Flavour Constraints on RR Sector

- * No Chargino Contributions
- * Only Neutralino contributions
- * Partial cancelations between the two contributions

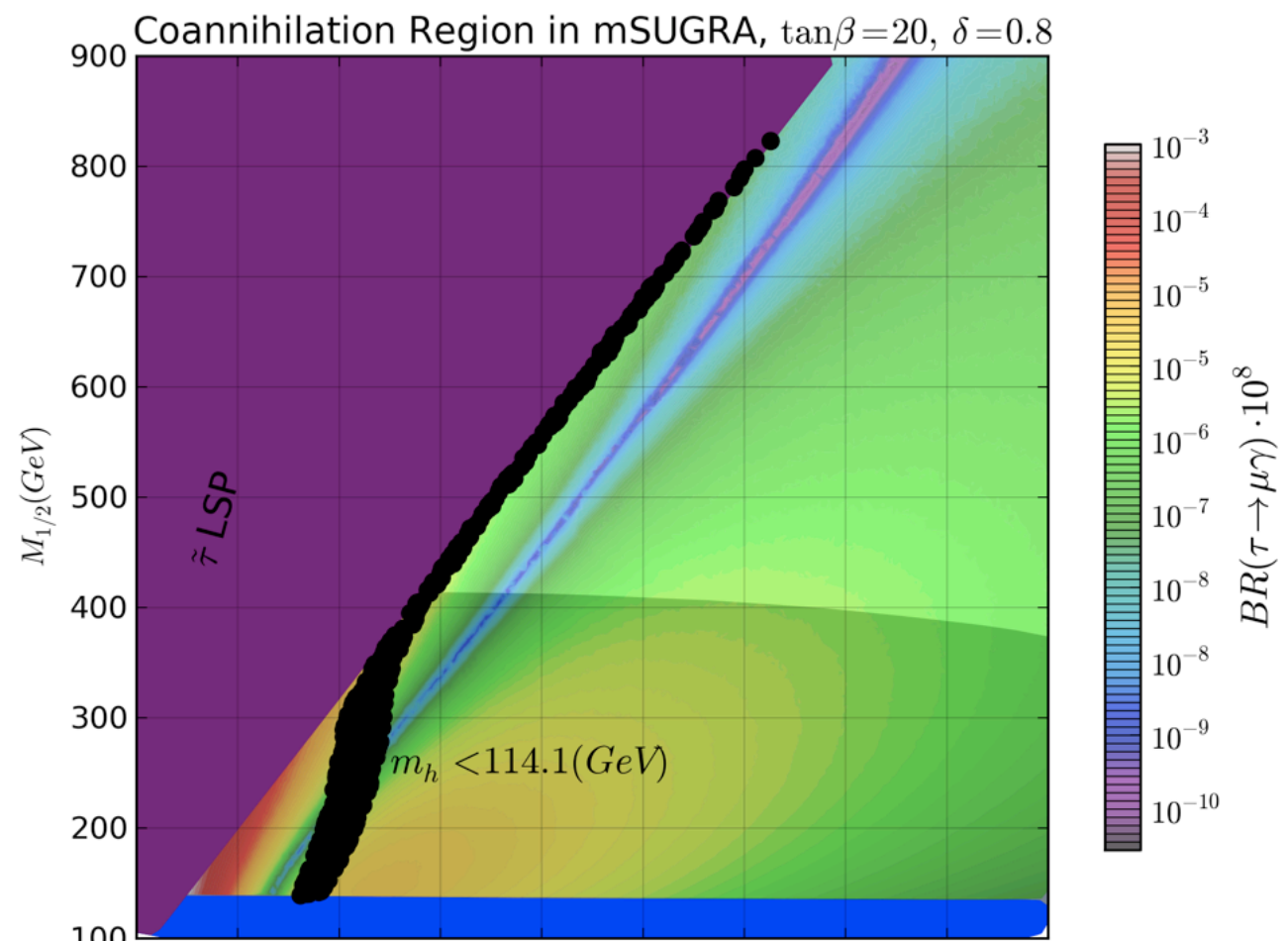
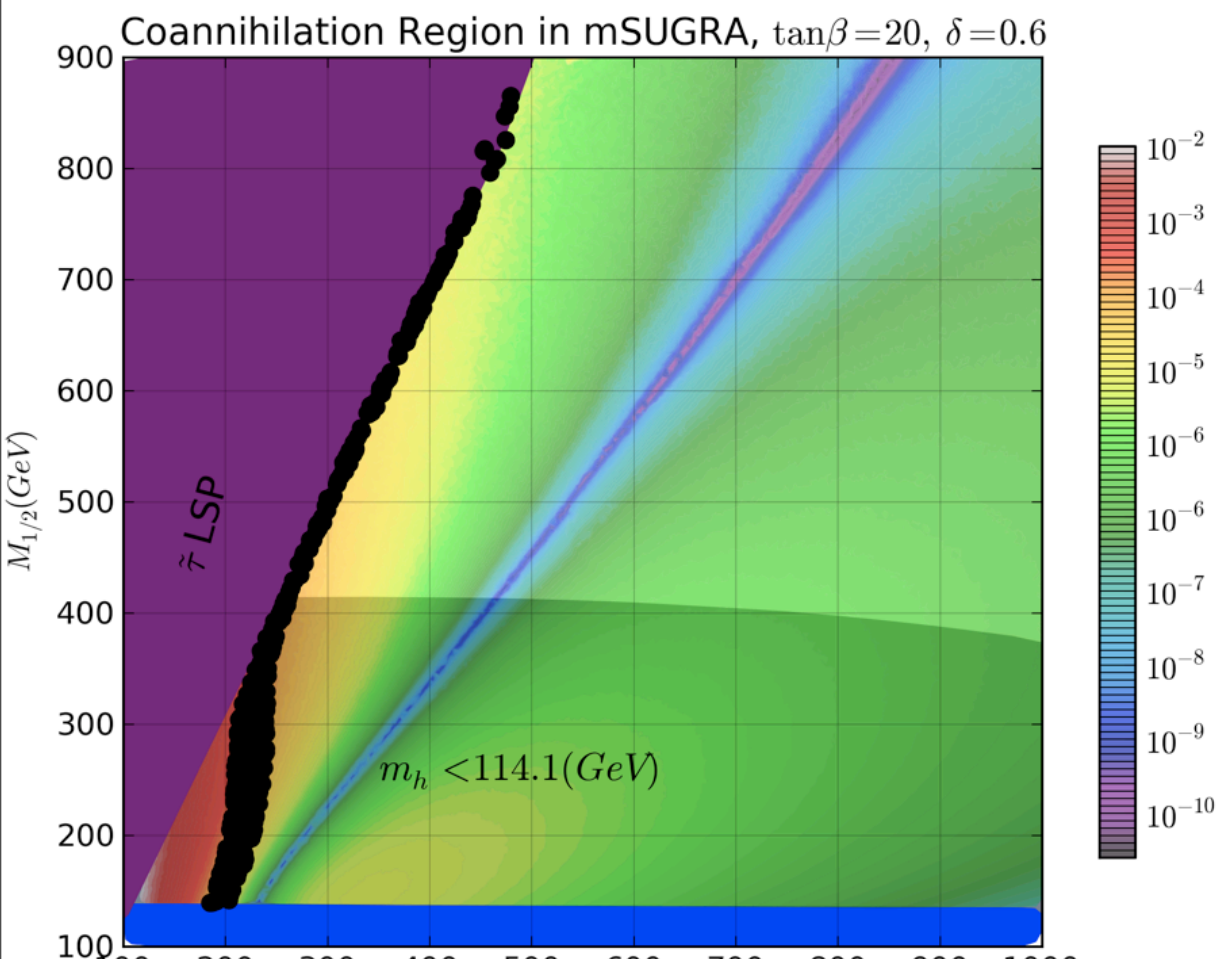
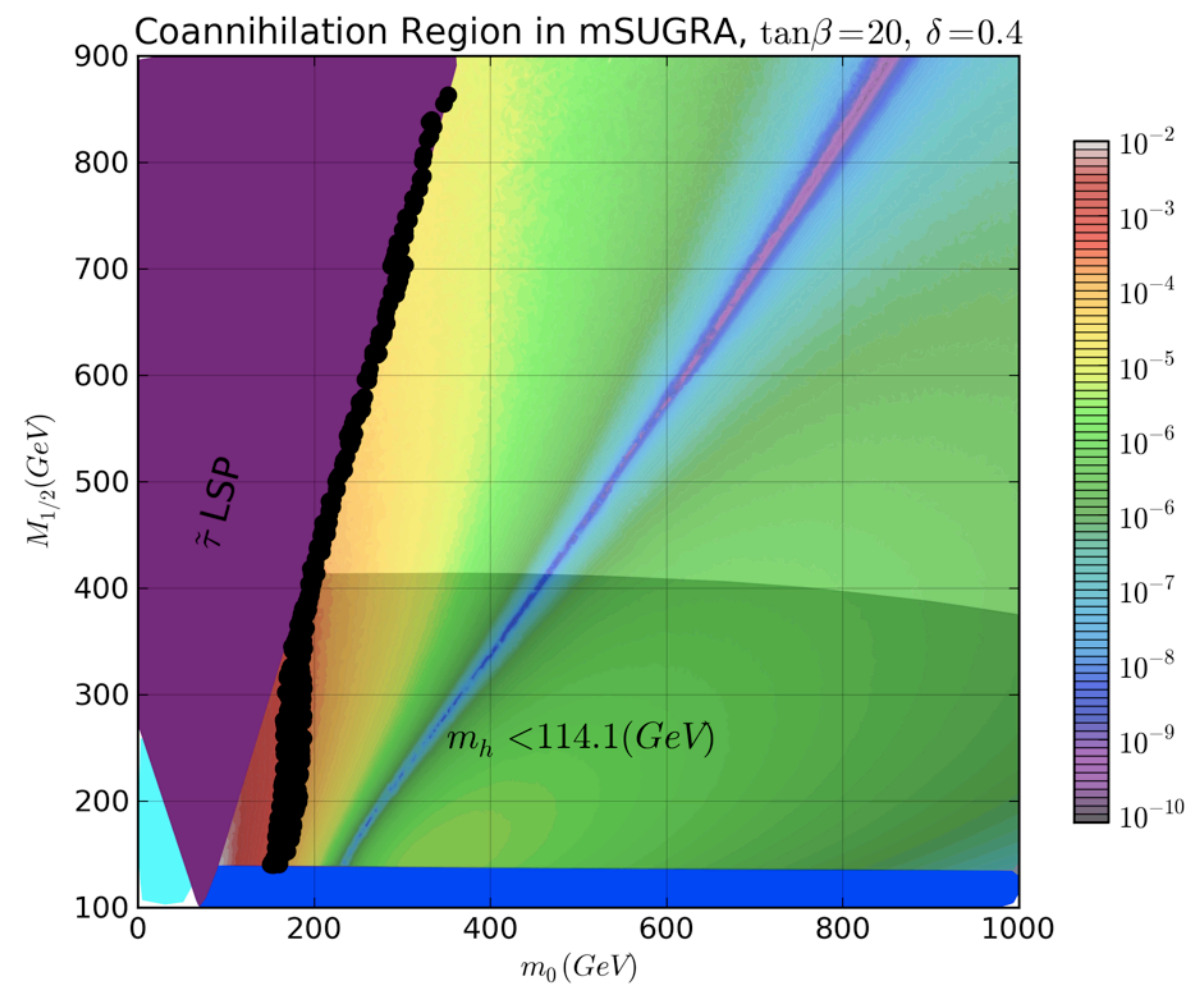
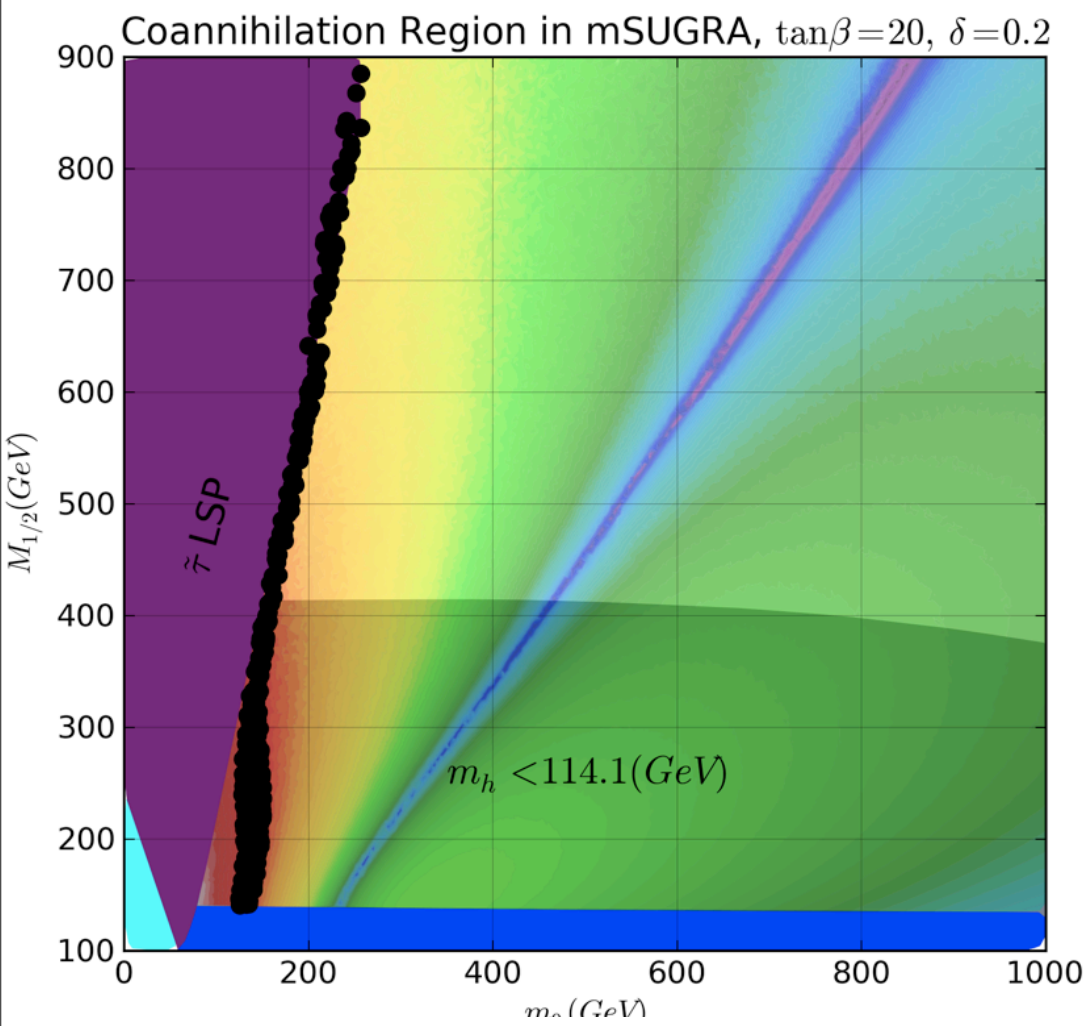
Hisano et.al 96
Hisano Nomura 99
Paride 06
Ciuchini et.al 07



From Masina and Savoy, NPB '02



There is no overlap between both the regions in mSUGRA



Relaxing the universality

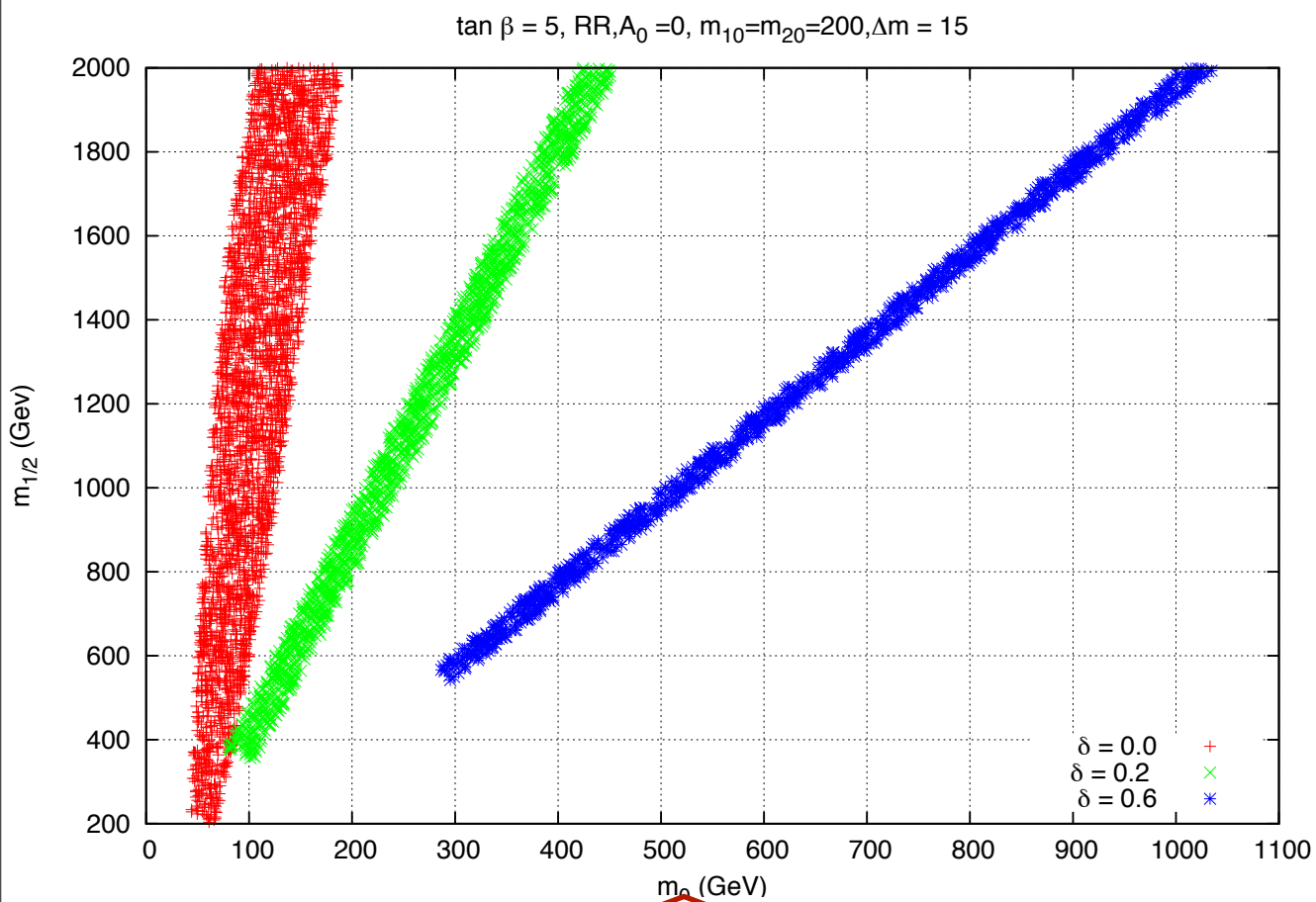
- * In GUT models like SO(10), all the matter sits in a single multiplet. It is likely that if soft masses are non-universal at origin, it would manifest in the higgs.

Consider NUHM with $m_{H_1^0}^2 = m_{H_2^0}^2$

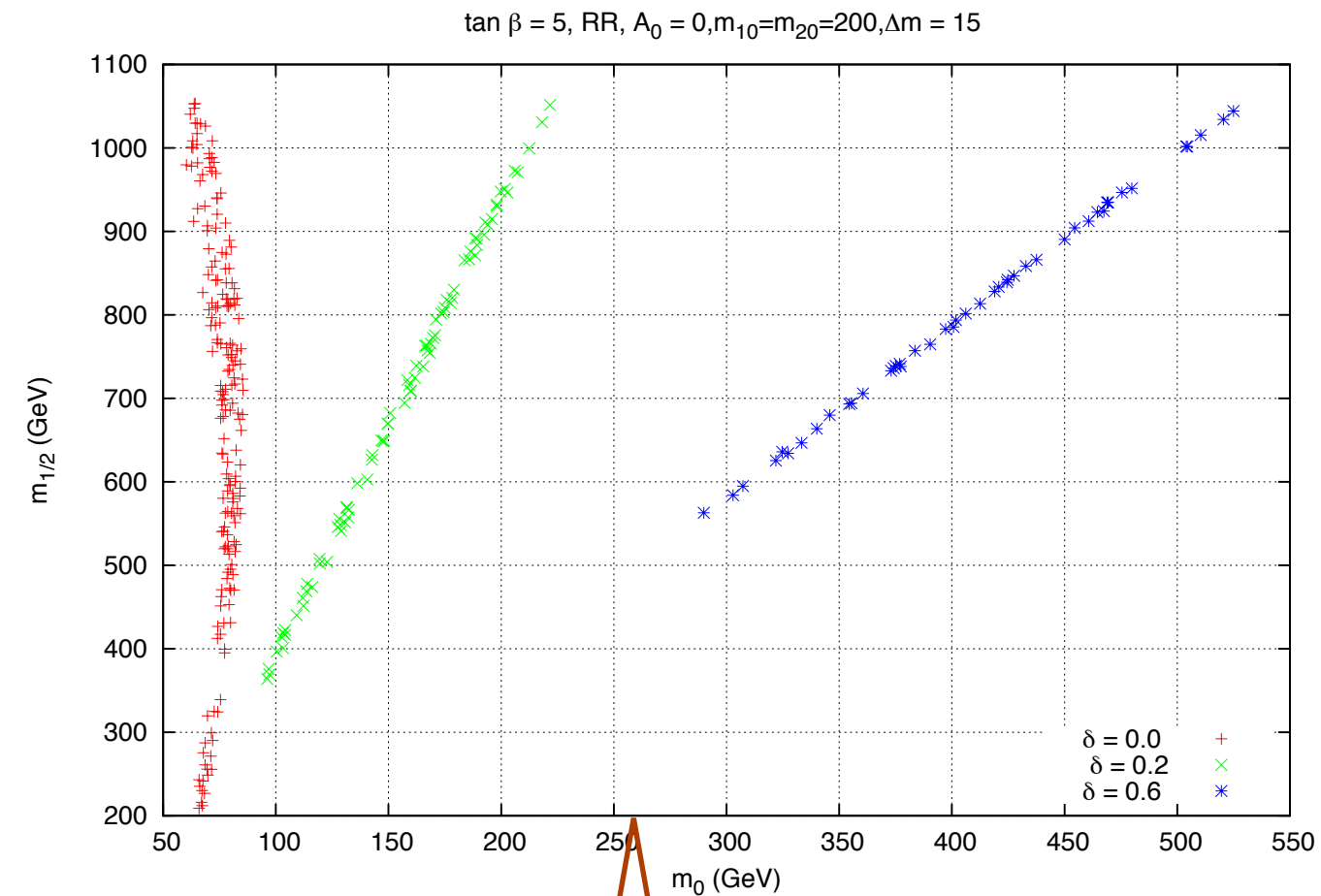
Overlap between regions where
cancelations in flavour violation
and regions with co-annihilations
is now possible !

μ becomes a free parameter

NUHM model

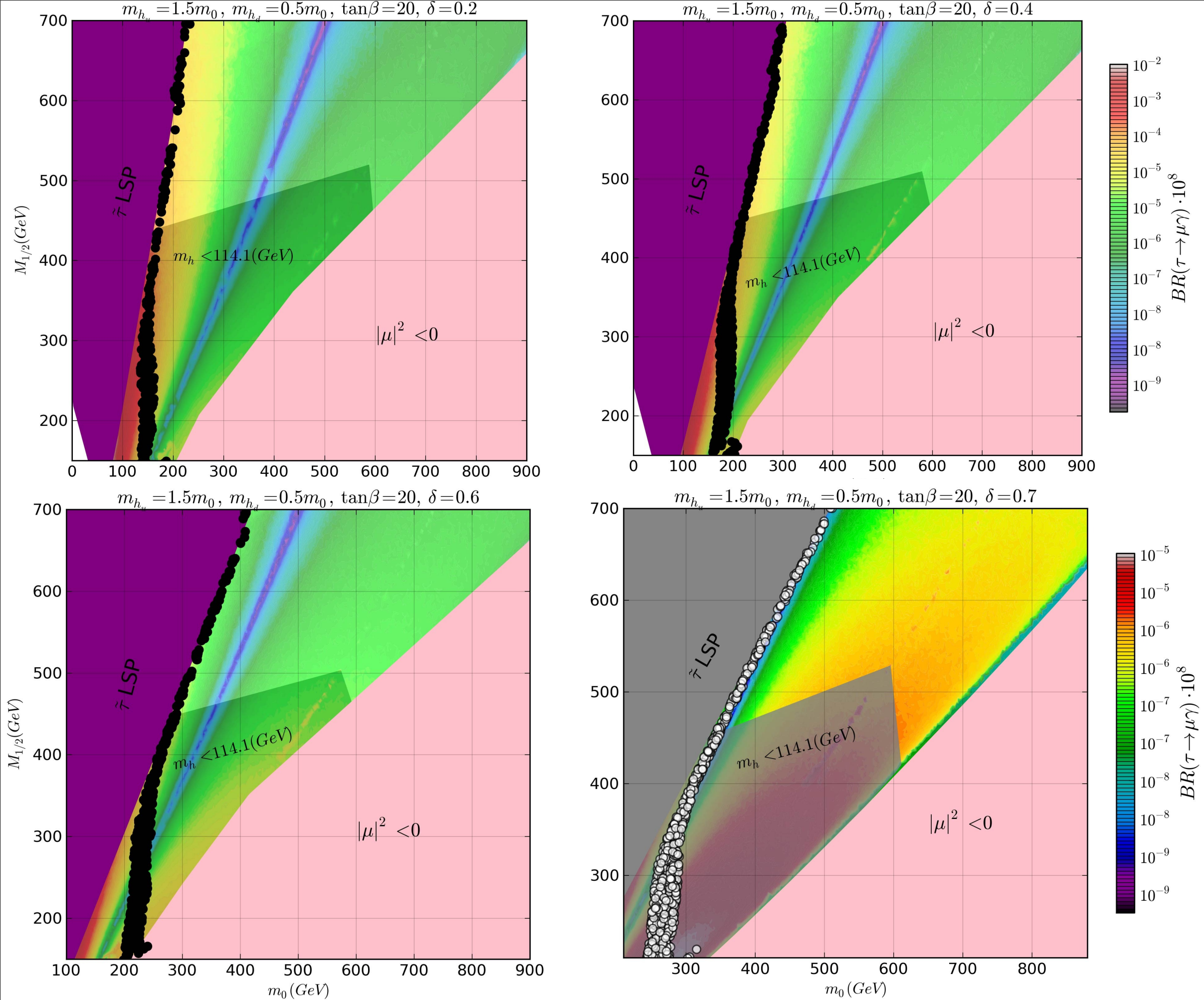


Imposing all constraints
except flavour violation

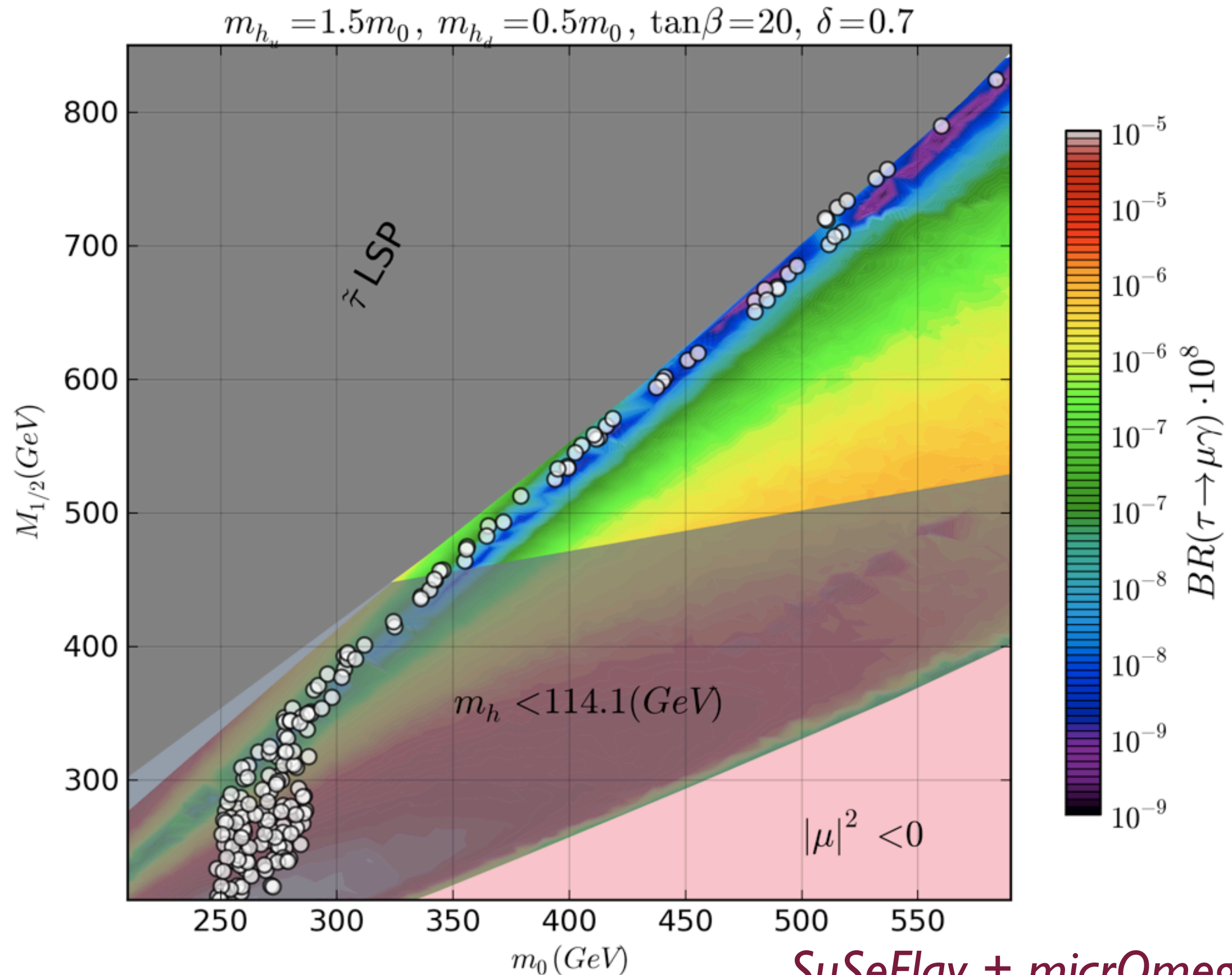


Including flavour
violating constraints

$$\text{Br}(\tau \rightarrow \mu + \gamma) < 2 \times 10^{-8}$$

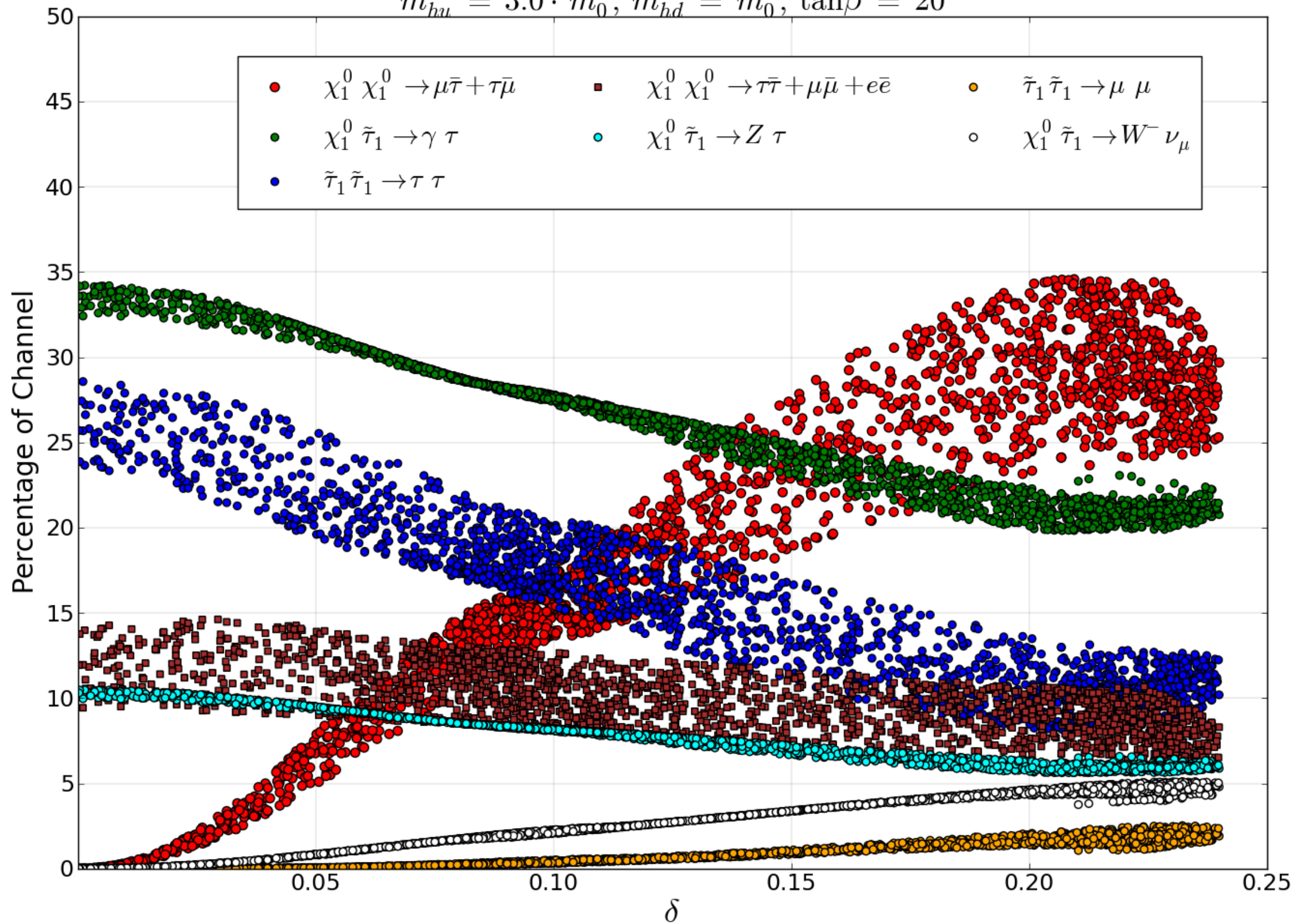


Flavoured Coannihilations in NUHM



SuSeFlav + micrOmegas

$$m_{hu} = 3.0 \cdot m_0, m_{hd} = m_0, \tan\beta = 20$$



Regions with smaller delta also
can be found

Summary of Flavour effects

Large flavour violation if present in the RR sector can modify significantly the co-annihilation region both in spectrum and relic density calculations

Extensions of mSUGRA like NUHM have interesting regions with *flavoured co-annihilations*

Low $\tan\beta$ is more preferable compared to high $\tan\beta$ as the amplitudes increase as $\tan\beta$

A. Datta et. al
Gomez et. al

We expect the regions to give interesting lepton flavour violating signals at LHC/ILC